

# Spaceborne Lidar Calibration From Cirrus and Molecular Backscatter Returns

J. A. Reagan, *Fellow, IEEE*, X. Wang, and M. T. Osborn

**Abstract**—In order to make optimal quantitative use of multiwavelength spaceborne lidar data, it is essential that the lidar be well calibrated. Due to system gain/efficiency changes that can be expected to occur during the course of a shuttle or satellite mission, it is essential to employ a calibration approach that can be implemented on-orbit, preferably repeatable at least a few times per orbit. For wavelengths less than about 550 nm, *in situ* calibration can be accomplished via normalization to high-altitude nearly molecular scattering regions. However, for longer wavelengths beyond about 800 nm, particularly the popular Nd: YAG fundamental wavelength at 1064 nm, the Rayleigh normalization approach becomes questionable due to both an inherently weaker signal and a stronger, variable, and somewhat unknown aerosol scattering contribution. For lidars operating at both longer and shorter wavelengths, a viable approach is to retrieve the longer wavelength calibrations ratioed to the shorter wavelength calibrations via comparisons of spectral backscatter from known/quantifiable scatterers. Cirrus clouds are good for this purpose because they occur at high altitudes with significant frequency and provide strong nearly spectrally flat backscatter. This paper presents both the molecular normalization and cirrus spectral backscatter ratio calibration approaches, including results obtained from case studies of lidar data collected during the LITE shuttle mission. Attention is focused on developing a simple autonomous approach applicable to satellite lidar missions such as Cloud-Aerosol Lidar Infrared Pathfinder Satellite Observations (CALIPSO) and the Geoscience Laser Altimeter System (GLAS).

**Index Terms**—Cirrus backscatter, Lidar calibration, spaceborne Lidar.

## I. INTRODUCTION

CLOUD-AEROSOL Lidar Infrared Pathfinder Satellite Observations (CALIPSO) (formally referred to as PICASSO-CENA [1]) is a recently approved mission within the National Aeronautic and Space Administration's (NASA) Earth System Science Pathfinder (ESSP) program that will obtain new measurements of clouds and aerosols to improve predictions and the impact of climate change. This mission is being developed as a partnership between NASA and the French space agency CNES and is planned for launch in early 2004. The CALIPSO payload consists of a two-wavelength (532 nm and 1064 nm) polarization-sensitive (for 532 nm) lidar and two coaligned passive instruments. Data from these three instruments will be used to measure the vertical distributions

of aerosols and clouds in the atmosphere, as well as optical and physical properties of aerosols and clouds, that influence the earth's radiation budget. The Geoscience Laser Altimeter System (GLAS) laser altimeter to be carried on ICESat (one of NASA's Earth Observing System satellites planned for launch at the end of 2002 or early 2003), which has the primary task of measuring ice-sheet topography, also features a two-wavelength atmospheric lidar [2]. Calibration of these satellite lidars is essential in order to achieve quantitative retrievals of aerosol and cloud properties.

Following the procedures employed for the Lidar In-Space Technology Experiment (LITE), *in situ* calibration of the 532-nm lidar channel can be accomplished via normalization to a high-altitude nearly molecular scattering region. However, the molecular backscatter is too weak to permit such calibration for the longer 1064-nm channel. An alternative approach investigated in [3] during the LITE shuttle mission [4] employed surface backscatter returns from selected land surfaces as standard targets to calibrate the LITE 1064-nm channel. This approach yielded calibration uncertainties in the 10% range, but only for limited well-characterized surfaces. During the course of a satellite mission, a more desirable calibration approach is one that can be implemented on-orbit, preferably repeatable at least a few times per orbit. With this objective, a new 1064-nm calibration method is proposed in this paper, specifically that the calibration of the 1064-nm channel relative to the 532-nm channel calibration be accomplished via comparisons of the 532- and 1064-nm backscatter signals from cirrus clouds. Due to the low signal for noncloud returns, it is difficult to remove the noncloud signal from the total signal. For the approach presented here, only strong cloud signal returns are selected. Consequently, the contamination by noncloud returns is minimized. This avoids uncertainty in subtracting a variable background level that must be done for weak cloud returns.

Examples demonstrating the 532-nm molecular normalization calibration approach are presented based on representative LITE data from orbits 24 and 34. Error assessments are included that support achievable calibration uncertainties within 5% by this approach. Cirrus cloud returns from LITE orbits 23, 24, and 27 have been analyzed to assess the feasibility of the cirrus cloud calibration approach. Results are given which indicate that calibration of the CALIPSO 1064-nm channel in terms of, or as a ratio to, the 532-nm calibration factor by using cirrus cloud returns appears quite feasible.

## II. PHYSICAL MODEL AND MATHEMATICAL DESCRIPTION

Calibration of a nadir/near-nadir viewing spaceborne lidar via normalization to high-altitude nearly molecular scattering

Manuscript received September 30, 2001; revised February 10, 2002. The work is supported by the National Aeronautics and Space Administration Langley Research Center under Contract NAS1-99102.

J. A. Reagan and X. Wang are with the Department of Electrical and Computer Engineering, University of Arizona, Tucson, AZ 85721 USA.

M. T. Osborn is with the Science Applications International Corp., NASA Langley Research Center, Hampton, VA 23681 USA.

Digital Object Identifier 10.1109/TGRS.2002.802464

regions is, in principle, very straightforward [3]. The lidar calibration factor or constant,  $C$ , appears in the lidar equation and normalized equation as follows (e.g., [5], [6]):

$$P(r) = \frac{E_0 C \beta(r) T^2(r)}{r^2} \quad (1)$$

$$X(r) = \frac{r^2 P(r)}{E_0} \quad (2)$$

where

- $E_0$  transmitted laser pulse energy;
- $r$  range or distance from the lidar to the point of scattering;
- $P(r)$  instantaneous lidar signal from range  $r$ ;
- $\beta(r)$  atmospheric backscatter coefficient ( $m^{-1} sr^{-1}$ );
- $T(r)$  atmospheric transmittance through range  $r$ .

The calibration constant for 532 nm may be extracted from the lidar signal  $P(r_c)$  obtained at a reference calibration range  $r_c$  by

$$C_{532} = \frac{r_c^2 P(r_c)}{R(r_c) \beta_m(r_c) T^2(r_c) E_0} = \frac{X(r_c)}{R(r_c) \beta_m(r_c) T^2(r_c)} \quad (3)$$

where, in addition to the terms defined above,

- $\beta_m(r_c)$  molecular (Rayleigh) atmospheric backscatter (for the lidar wavelength) at range  $r_c$  for 532 nm;
- $R(r_c)$   $\beta(r_c)/\beta_m(r_c)$ ; total-to-molecular backscattering mixing ratio at range  $r_c$

For  $r_c$  selected to be around 30 km above ground,  $R(r_c) \approx 1$  and  $T^2(r_c) \approx 0.99$  are very good approximations for 532 nm, enabling accurate retrievals of  $C_{532}$ , providing that the signal uncertainty is sufficiently small and that  $\beta_m(r_c)$  can be accurately computed (driven by how accurately the air density can be determined).

The total Rayleigh-scattering cross section per molecule  $\sigma$  is given by [7], [8]

$$\sigma(\lambda) = \frac{24\pi^3(n_s^2 - 1)^2}{\lambda^4 N_s^2(n_s^2 + 2)^2} \left( \frac{6 + 3\rho_n}{6 - 7\rho_n} \right) \quad (4)$$

where  $\lambda$  is the wavelength (in centimeters);  $n_s$  is the refractive index for standard air at  $\lambda$ ;  $N_s$  is the molecular number density for standard air; and  $\rho_n$  is the depolarization factor, a term that accounts for the anisotropy of the air molecule and which varies with wavelength. The total molecular volume backscattering coefficient  $\beta_{m,\lambda}(z)$  is given by the product of the total Rayleigh cross section per molecule  $\sigma$  as defined in (4), the molecular number density  $N(z)$ , for a given pressure and temperature at altitude  $z$ , and the Rayleigh extinction-to-backscatter ratio  $S_R$  (e.g., see [6] and [9])

$$\beta_{m,\lambda}(z) = S_R^{-1} N(z) \sigma(\lambda) \quad (5)$$

where  $S_R = 8\pi/3$ .

The Rayleigh-per-molecule cross section can be accurately determined within a few tenths of a percent [7], leaving  $N(z)$  as the primary source of uncertainty in determining  $\beta_{m,\lambda}(z)$ . Using ancillary meteorological data along the satellite track, it is estimated that  $\beta_m(r_c)$  can be determined within about  $\pm 3\%$

uncertainty. For  $z$  near 30 km, this corresponds to knowing temperature within  $\sim \pm 5$  K and pressure within  $\sim \pm 2.5$  mb, which is about the level of uncertainty associated with just assuming a latitudinal, seasonal standard atmospheric model [5], [8]. Using pressure and temperature fields derived from weather network measurements and assimilation models, coupled with averaging the  $C_{532}$  retrievals over significant horizontal extents ( $\sim$ several hundred kilometers), it is anticipated that the  $N(z)$  uncertainty can be further reduced. However, for the estimates presented here, it is assumed that  $\beta_m(r_c)$  is not known better than  $\pm 3\%$ .

Given  $C_{532}$  determined by the molecular calibration approach, cirrus clouds offer good candidate targets by which the calibration ratio  $C_{1064}/C_{532}$  can be estimated from the ratio of the normalized returns for the two wavelengths. This is feasible because, to first order, the backscatter and extinction from cirrus should be nearly the same for both wavelengths. Also, as cirrus clouds occur at high altitudes, corrections for 1064/532 spectral transmittance differences between the satellite and the cloud top are relatively small and predictable.

The normalized cloud return  $X_c(r)$ , which is defined as the total normalized return minus the noncloud background normalized return, is given approximately by

$$X_c(r) = C_{1064} T_{ct}^2 \beta_c(r) T_c^2(r) \quad (6)$$

where

- $C_{1064}$  lidar calibration factor;
- $T_{ct}^2$  round-trip transmittance to cloud top at range  $r_{ct}$ ;
- $\beta_c(r)$  cloud backscatter for  $r > r_{ct}$ ;
- $T_c^2(r)$  cloud round-trip transmittance from  $r_{ct}$  to  $r > r_{ct}$ .

As cirrus particles are typically sufficiently large relative to the 532- and 1064-nm wavelengths for the geometrical optics limit to reasonably apply, there should be no wavelength dependence in the extinction and backscatter for these two wavelengths [10], [11], save any small refractive index difference effects. In addition, cirrus spectral optical depths measured by sunphotometers are observed to be quite spectrally flat over the  $\sim 500$ -nm to the  $\sim 1000$ -nm range [12]. In fact, screening for spectral flatness in optical depth is a technique often employed to detect the presence of very thin cirrus contamination in derived aerosol optical depths. Lidar data, though sparse, supports  $\beta_{c,532} = \beta_{c,1064}$  within measurement/calibration uncertainties (albeit, these uncertainties are typically somewhat large,  $\sim 10\%$ – $20\%$ ). Preliminary analysis of lidar observations made during the SAFARI-2000 campaign with the NASA/GSFC Cloud Physics Lidar, deployed on a NASA ER-2, reveals that  $\beta_{c,1064}/\beta_{c,532}$  is quite constant, not highly variable, and close to unity [13]. Finally, for spaceborne lidar to cloud geometry and even a very small receiver field of view ( $\sim 100 \mu\text{rad}$ ), multiple scattering effects should be effectively the same for both wavelengths. All of this supports assuming  $\beta_c T_c^2$  to be spectrally flat.

Assuming  $\beta_c T_c^2$  is the same for 532 nm and 1064 nm, the ratio of  $X_c$  for the two wavelengths at any  $r$  within the cloud will be

$$\frac{X_{c,1064}}{X_{c,532}} = \frac{C_{1064}}{C_{532}} \times \frac{T_{ct,1064}^2}{T_{ct,532}^2} \quad (7)$$

yielding

$$\frac{C_{1064}}{C_{532}} = \frac{X_{c1064}}{X_{c532}} \times \frac{T_{ct,532}^2}{T_{ct,1064}^2}. \quad (8)$$

In order to minimize the noncloud return, a threshold is selected to screen out the weak backscatter signals. By setting a high threshold, the noncloud background can be ignored, allowing  $X_c$  to be approximated by  $X$ , the total signal including the noncloud background. Thus, the calibration ratio is given approximately by

$$\frac{C_{1064}}{C_{532}} \approx \frac{X_{1064}}{X_{532}} \times \frac{T_{ct,532}^2}{T_{ct,1064}^2}. \quad (9)$$

The last term is approximately 0.9 at about 12 km above ground and can be estimated from models of aerosol extinction and ozone concentration and an atmospheric density profile [8].

A threshold signal  $X_t$  for determining strong cloud returns may be determined by computing the 532-nm normalized signal that is equivalent to a scattering ratio of  $R_t$ .

Expressing  $X_t(r)$  in terms of altitude above ground,  $z$ , which is normally how molecular and aerosol scattering coefficients and cloud positions are height referenced,  $X_t(z)$  may be expressed by

$$X_t(z) = C_{532} \times R_t \times \beta_{m,532}(z) \times T_{532}^2(z_L - z) \quad (10)$$

where  $R_t = (\beta_a(z) + \beta_m(z))/(\beta_m(z))$ ;  $z$  is the observation altitude above ground; and  $z_L$  is the lidar height above ground (i.e.,  $r$  and  $z$  are related by  $r = z_L - z$ ).

Setting  $R_t$  to a large value, on the order of 50, insures  $X_c \approx X$ . Once again, for the heights of cirrus clouds, the transmission term can be modeled accurately and is close to unity. This threshold is applied only to the 532-nm signals for determining cloud segments of sufficient signal intensity to be used for retrieving the calibration ratio.

### III. RESULTS AND DISCUSSION

#### A. Retrieval Results for the 532-nm Calibration Factor

Extracting absolute estimates of  $C_{532}$ , as can be seen from (3), requires that  $r_c$ ,  $E_0$ ,  $R_{532}(r_c)$ ,  $\beta_{m,532}(r_c)$ , and  $T_{532}^2(r_c)$  all be known/specified. It is assumed that this can be done with essentially no significant error (i.e., less than  $\sim 0.5\%$  error) for  $r_c$  and  $E_0$  (a relative energy normalization is all that is actually required), and with manageably small error for  $\beta_{m,532}(r_c)$ , using temperature and pressure meteorological data incorporated in the CALIPSO database. As discussed earlier, using molecular number densities calculated from pressure and temperature fields derived from assimilated network measurements should yield estimates of  $\beta_{m,532}(r_c)$ , easily within  $\pm 3\%$ . Also, for  $r_c$  selected in the mid to upper stratosphere,  $T_{532}^2(r_c)$  will generally be within about 1% of unity, but must still be specified to extract  $C_{532}$ .

In Fig. 1, retrievals of  $C_{532}$  from orbit 34 LITE data are presented for 4-km vertical averaging over the altitude range  $z = 30$  to 34 km and 200-shot ( $\sim 150$  km) minimum horizontal block averaging, for various total horizontal extents of

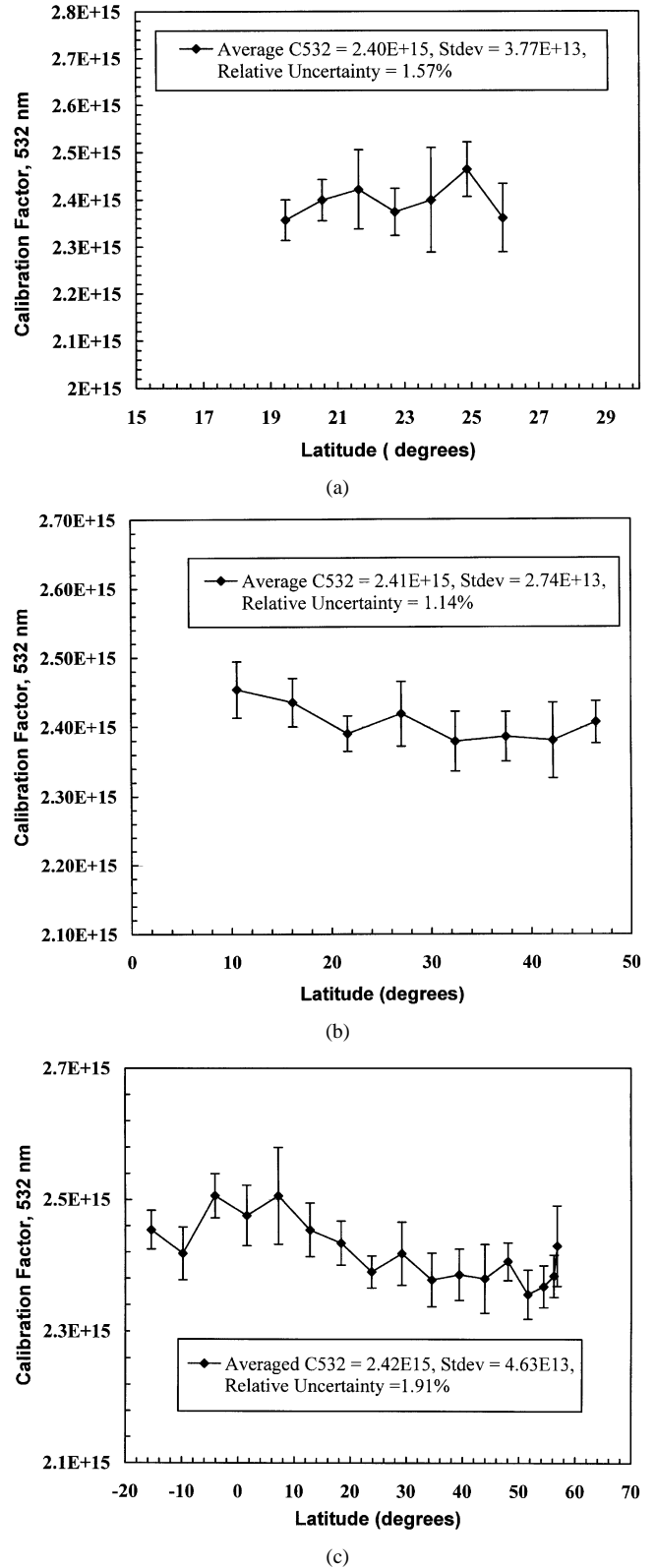


Fig. 1. Retrieved 532-nm calibration constant for portions of LITE orbit 34. (a) Based on a 200-shot average spanning about a 1000-km horizontal extent. (b) Based on a 1000-shot horizontal average over about half the nighttime portion of orbit 34. (c) Based on a 1000-shot horizontal average for the entire nighttime portion of orbit 34.

$\sim 1000$  km up to half an orbit. In Fig. 1(a),  $C_{532}$  is calculated based on a 200-shot average spanning about a 1000-km total horizontal extent (during the nighttime portion). Fig. 1(b)

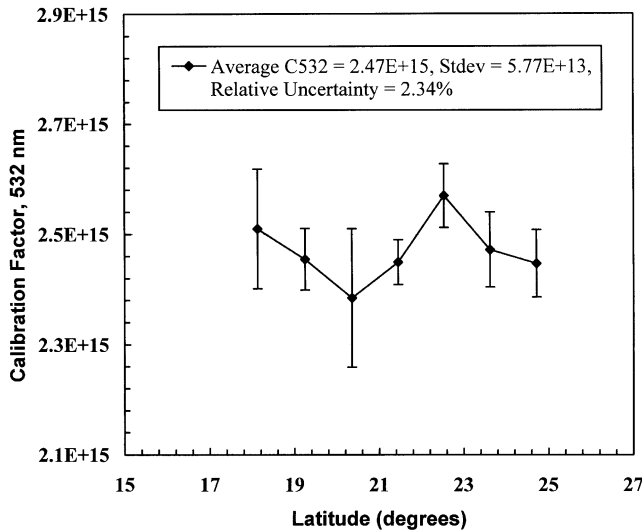


Fig. 2. Retrieved 532-nm calibration constant for a portion of LITE orbit 24 based on a 200-shot average spanning about a 1000-km horizontal extent.

is based on a 1000-shot average, spanning about half the nighttime portion of the orbit, and in Fig. 1(c), the 1000-shot average extends over the entire nighttime portion of the orbit. Each data point in Fig. 1(a)–(c) includes  $\pm 1$  standard deviation error bar for the number of shots averaged for the point (either 200 or 1000 shots). These values are averaged for the number of points in each figure, yielding the average and standard derivation values cited in the figures. The standard deviation of the 200-shot average points is mainly determined by the uncertainty in the normalized signal  $X$  and is in the range of 3%.

From Fig. 1, it can be seen that  $C_{532}$  for orbit 34 is around  $2.41 \times 10^{15}$ , and the relative uncertainty of  $C_{532}$  is less than 2%. Here it is important to note that these figures depict the effects of signal uncertainty and relative spatial variability in  $\beta_{m,532}(r_c)$ ,  $R_{532}(r_c)$ , and  $T_{532}^2(r_c)$ , but not the absolute errors in these latter three factors. What the figures demonstrate is that averaging over horizontal extents of  $\sim 1000$ – $10\,000$  km should still yield averages with uncertainties within about 2% insofar as these horizontal inhomogeneity effects are concerned. The results shown for orbit 34 are typical of other orbits [14], another example of which is shown in Fig. 2 for a segment of orbit 24, yielding a  $C_{532}$  within about 2% of the orbit 34 results. Some change in  $C_{532}$  between several orbits was observed during the LITE mission [14], presumably reflecting instrument changes and drift. Hence, a repeatable on-orbit calibration approach is needed to overcome these effects.

Applying standard error propagation analysis to the  $C_{532}$  retrieval equation (2) yields

$$\left[ \frac{\delta C_{532}}{C_{532}} \right]^2 = \left[ \frac{\delta X_{532}(r_c)}{X_{532}(r_c)} \right]^2 + \left[ \frac{\delta R_{532}(r_c)}{R_{532}(r_c)} \right]^2 + \left[ \frac{\delta \beta_{m,532}(r_c)}{\beta_{m,532}(r_c)} \right]^2 + \left[ \frac{\delta T_{532}^2(r_c)}{T_{532}^2(r_c)} \right]^2. \quad (11)$$

Assuming uncertainties that should reasonably apply in the stratosphere for  $r_c$  in the range of 30–34 km above ground, as listed in Table I, yields an uncertainty in  $C_{532}$  of about 4.4% as also given in Table I.

TABLE I  
THE UNCERTAINTY ESTIMATES OF THE 532-nm  
CALIBRATION CONSTANT AND RELATED PARAMETERS

$\frac{\delta C_{532}}{C_{532}}$	$\frac{\delta X_{532}}{X_{532}}$	$\frac{\delta R_{532}(r_c)}{R_{532}(r_c)}$	$\frac{\delta \beta_{m,532}(r_c)}{\beta_{m,532}(r_c)}$	$\frac{\delta T_{532}^2(r_c)}{T_{532}^2(r_c)}$
0.044	0.03	0.01	0.03	0.005

TABLE II  
SUMMARY OF CALIPSO LIDAR PARAMETERS

Orbit Height	705 Km
Laser	Nd: YAG with $\sim 110$ mJ at 532 and 1064 nm
Polarization	532 nm has parallel and perpendicular channels
Laser Repetition Rate	20.25 Hz
Receiver Telescope Diameter	1.0 m
Vertical Resolution	30m - 300m, depending on altitude and wavelength
Horizontal Resolution	1/3 km - 5 km, depending on altitude

CALIPSO simulations based on specified lidar system parameters (summary of key parameters listed in Table II) predict that the shot-noise-limited uncertainty in  $X(r)$  for a height of 30 km above ground and 4-km vertical averaging should be within about  $\pm 3\%$  for horizontal averaging of not more than  $\sim 1000$  km. As shown in Figs. 1 and 2, results from LITE demonstrate that horizontal averaging over 1000 km and more without significant horizontal inhomogeneity biases is quite feasible. Hence, the molecular normalization approach should enable on-orbit calibrations of  $C_{532}$  for CALIPSO with uncertainties within  $\pm 5\%$ .

#### B. Retrieval Results for the $C_{1064}/C_{532}$ Calibration Ratio

The  $C_{1064}/C_{532}$  calibration ratio retrieval approach outlined in Section II was applied using LITE data for the nighttime portions of orbits 23, 24, and 27. The search for cloud returns was restricted to the altitude region from 8–17 km above ground. This helps eliminate noncirrus cloud returns and facilitates the modeling of transmission terms in the retrieval equations. The screening threshold level  $R_t$  was set to 50, insuring strong cloud returns, permitting minimal horizontal averaging on only 10 shots ( $\sim 7$  km horizontal extent) to still yield a signal uncertainty typically within about  $\pm 2\%$ . Also, only the highest altitude cloud with a thickness (as determined by the threshold signal) of at least 180 m was used for each 10-shot average. Fig. 3 shows the relationship between the retrieved calibration ratio  $C_{1064}/C_{532}$  and the latitudes spanned by an orbit, corresponding to orbits 23, 24, and 27, respectively. For orbit 23, 61 cloud profiles (each set being a 10-shot average) met the signal selection requirements specified for the retrievals. For orbits 24 and 27, 29 and 111 cloud profile sets, respectively, were used for retrievals. There appears to be slightly greater scatter in the points for higher latitudes. This may be due to some mixed-phase/water clouds being present at higher latitudes within the altitude screening range (8–17 km) that was used to select cirrus

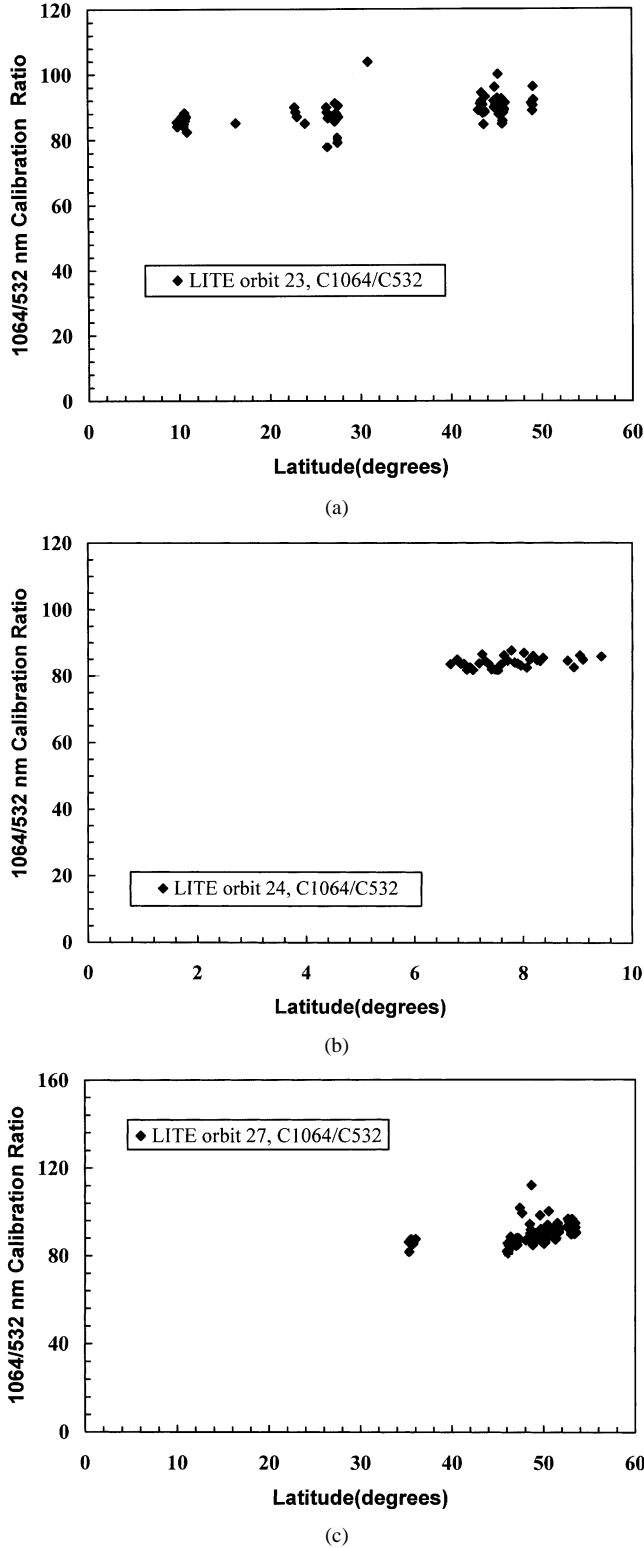


Fig. 3. Retrieved 1064/532 calibration ratios from selected cirrus cloud returns for (a) orbit 23, (b) orbit 24, and (c) orbit 27.

clouds. However, this variability does not appear to be statistically significant. The numerical results (means and standard deviations) for the  $C_{1064}/C_{532}$  calibration ratio retrievals are listed in Table III.

From Table III, it can be seen that the calibration ratio for the three nearly adjacent LITE orbits 23, 24, and 27 is very con-

TABLE III  
CALIBRATION RATIO  $C_{1064}/C_{532}$  FOR LITE ORBITS 23, 24, AND 27

Orbit	Calibration Ratio $C_{1064}/C_{532}$	Number of Cloud Profiles
23	$88.1 \pm 2.3$	61
24	$85.3 \pm 1.5$	29
27	$88.5 \pm 2.6$	111

tent with small standard deviations within  $\pm 3\%$ . In addition, the calibration ratio of these orbits agrees with the calibration ratio obtained from LITE ground reflections over Edwards dry lake bed (during orbit 24) within uncertainties for the ground reflection calibration retrieval [3]. The calibration ratio  $C_{1064}/C_{532}$  obtained at Edwards Air Force Base, CA was about  $79 \pm 8$ .

### C. Retrieval Uncertainty Estimates for the Cirrus Cloud $C_{1064}/C_{532}$ Calibration Approach

Assuming the current parameters for CALIPSO, the expected signal-to-noise ratio (SNR) for single-shot returns at about 11 km above ground from cirrus clouds with backscatter 50 times greater than 532-nm molecular backscatter, for a 180-m vertical range bin, is  $\text{SNR} \approx 5$  for both 532 nm and 1064 nm. Averaging over 5 km (15 shots) horizontally should yield an SNR of about 20, or a signal uncertainty ( $1/\text{SNR}$ ) of about 5%, which reduces the signal uncertainty to a small enough level to apply fairly stringent statistical discriminants to reject “contaminated” or “atypical” horizontal cirrus cloud segments. Averaging over additional horizontal segments and/or vertical range bins should further reduce signal uncertainties as demonstrated by the LITE examples.

Including the possibility of some difference between  $\beta T^2$  for 532 nm and 1064 nm within the cloud ( $r > r_{ct}$ ), the ratio of the calibration constants may be expressed as

$$\frac{C_{1064}}{C_{532}} = \frac{X_{1064}}{X_{532}} \times \frac{T_{ct,532}^2}{T_{ct,1064}^2} \times \chi' \quad (12)$$

where

$$\chi' = \frac{\beta_{532}}{\beta_{1064}} \times \frac{T_{r>ct,532}^2}{T_{r>ct,1064}^2} \quad (13)$$

and

$\beta_\lambda$  total cloud and noncloud backscatter;  
 $T_{r>ct,\lambda}^2$  total two-way transmittance from  $r$  to  $r > r_{ct}$ .

Applying standard error propagation analysis to (12), the relative uncertainty in  $C_{1064}$  can be expressed as

$$\left[ \frac{\delta C_{1064}}{C_{1064}} \right]^2 = \left[ \frac{\delta X_{1064}}{X_{1064}} \right]^2 + \left[ \frac{\delta X_{532}}{X_{532}} \right]^2 + \left[ \frac{\delta T_{ct,532}^2}{T_{ct,532}^2} \right]^2 + \left[ \frac{\delta T_{ct,1064}^2}{T_{ct,1064}^2} \right]^2 + \left[ \frac{\delta C_{532}}{C_{532}} \right]^2 + \left[ \frac{\delta \chi'}{\chi'} \right]^2. \quad (14)$$

Conservative estimates for the error terms in (14) are given in Table IV, along with the resulting error in  $C_{1064}$ . Hence, it appears quite feasible to determine  $C_{1064}$  within  $\pm 10\%$  using the

TABLE IV  
UNCERTAINTY ESTIMATES OF THE 1064-nm CALIBRATION CONSTANT  
AND RELATED PARAMETERS

$\frac{\delta C_{1064}}{C_{1064}}$	$\frac{\delta X_{1064}}{X_{1064}}$	$\frac{\delta T_{cr,532}^2}{T_{cr,532}^2}$	$\frac{\delta X_{532}}{X_{532}}$	$\frac{\delta T_{cr,1064}^2}{T_{cr,1064}^2}$	$\frac{\delta C_{532}}{C_{532}}$	$\frac{\delta \chi'}{\chi'}$
0.097	0.05	0.02	0.05	0.002	0.05	0.04

cirrus cloud ratio approach. Given the ratio results from the LITE data analysis, it is anticipated that the  $C_{1064}$  calibration uncertainty that can be achieved for CALIPSO may more likely be less than  $\pm 10\%$ .

#### IV. SUMMARY AND CONCLUSION

Techniques have been presented for on-orbit calibration of spaceborne lidar applicable for both shorter visible wavelengths, such as 532 nm, and longer near-infrared wavelengths, such as 1064 nm. The approach for shorter wavelengths employs molecular (Rayleigh) scattering normalization at high altitudes ( $\sim 30$  km), while the approach for longer wavelengths determines the ratio of the longer-to-shorter wavelength calibration factors from selected cirrus cloud backscatter returns. These approaches have been demonstrated using example data from the LITE shuttle mission, including assessment of uncertainties.

Results presented here show that calibration of the CALIPSO 532-nm channel by molecular normalization, analogous to what was done for the LITE shuttle mission, is quite feasible and should yield calibrations with uncertainties of  $\pm 5\%$  or less. Calibration of the CALIPSO 1064-nm channel in terms of, or as a ratio to, the 532-nm calibration factor by using cirrus cloud returns also appears quite feasible. The accuracy with which this can be achieved should conservatively be  $\pm 10\%$  or less by selecting strong signal returns from cirrus clouds.

#### ACKNOWLEDGMENT

Information and assistance provided by other members of the NASA Langley Research Center are gratefully acknowledged.

#### REFERENCES

- [1] J. A. Reagan and D. Winker, "PICASSO-CENA: Combined active-passive sensing from space," in *Proc. IGARSS*, Hamburg, Germany, June-July 28-2, 1999, pp. 240-242.

- [2] J. D. Spinhirne and S. P. Palm, "Space based atmospheric measurements by GLAS," in *Selected Papers from 18th ILRC*, A. Ansmann, R. Neuber, P. Rairoux, and U. Wandinger, Eds. Berlin, Germany: Springer-Verlag, 1997, pp. 213-216.
- [3] J. A. Reagan, H. Liu, and T. W. Cooley, "LITE surface returns: Assessment and application, advances in atmospheric remote sensing with lidar," in *Selected Papers from 18th ILRC*, A. Ansmann, R. Neuber, P. Rairoux, and U. Wandinger, Eds. Berlin, Germany: Springer-Verlag, 1997, pp. 177-180.
- [4] D. M. Winker, R. H. Crouch, and M. P. McCormick, "An overview of LITE: NASA's lidar in-space technology experiment," *Proc. IEEE*, vol. 84, pp. 164-180, Feb. 1996.
- [5] P. B. Russell, T. J. Swissler, and M. P. McCormick, "Methodology for error analysis and simulation of lidar aerosol measurements," *Appl. Opt.*, vol. 18, pp. 3783-3797, 1979.
- [6] J. A. Reagan, M. P. McCormick, and J. D. Spinhirne, "Lidar sensing of aerosols and clouds in the troposphere and stratosphere," *Proc. IEEE*, vol. 77, pp. 433-448, Mar. 1989.
- [7] A. Bucholtz, "Rayleigh-scattering calculations for the terrestrial atmosphere," *Appl. Opt.*, vol. 34, no. 15, May 1995.
- [8] *Handbook of Geophysics and the Space Environment*, National Technical Information Service, Springfield, VA.
- [9] F. G. Fernald, B. M. Herman, and J. A. Reagan, "Determination of aerosol height distributions by lidar," *J. Appl. Meteorol.*, vol. 11, pp. 482-489, 1972.
- [10] P. Yang and K. N. Liou, "Finite-difference time domain method for light scattering by small ice crystals in three-dimensional space," *J. Opt. Soc. Amer.*, vol. A 13, pp. 2072-2085, 1996.
- [11] —, "Geometric-optics-integral-equation method for light scattering by nonspherical ice crystals," *Appl. Opt.*, vol. 35, pp. 6568-6584, 1996.
- [12] R. F. Pueschel and J. M. Livingston, "Aerosol spectral optical depths: Jet fuel and forest fire smokes," *J. Geophys. Res.*, vol. 95, pp. 22 417-22 422, 1990.
- [13] D. L. Hlavka, Personal Communication for the CALIPSO Lidar Algorithm Peer Review, NASA Goddard Space Flight Center, Hampton, VA, Aug. 1-2, 2001.
- [14] M. T. Osborn, "Calibration of LITE data," in *Proc. 19th Int. Laser Radar Conf.*, Annapolis, MD, July 6-10, 1998, pp. 245-247.

**J. A. Reagan**, photograph and biography not available at the time of publication.

**X. Wang**, photograph and biography not available at the time of publication.

**M. T. Osborn**, photograph and biography not available at the time of publication.